



WCCM XI - ECCM V - ECFD VI
BARCELONA 2014

11th. World Congress on
Computational Mechanics

Convergence analysis of configurational forces for brittle cracks modeled through C^k -Generalized FEM

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Presentation topics

- Construction of continuous **Partition of Unity** at C^k -GFEM
- Construction of a smooth approximation subspace
- Elshebian mechanics as tool to post-processing of J-integral
- Quality assessment through **global** and **local measures**
- Concluding remarks



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Continuous partition of
unity with C^k -GFEM

Defining an
approximation
subspace

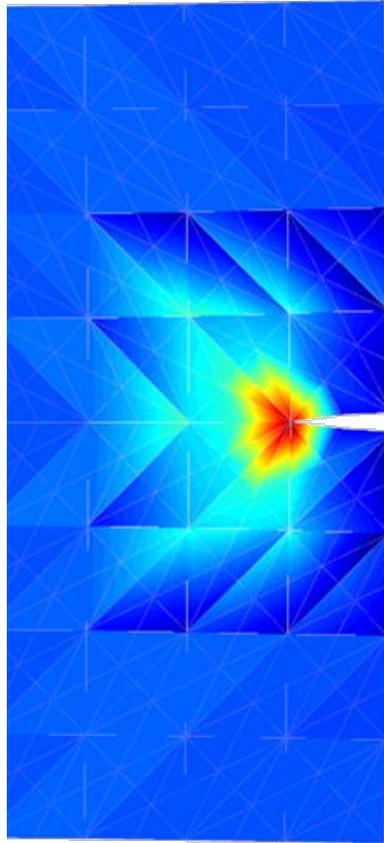
Quality assessment
through global
measures

Configurational forces
method

Quality assessment
through local
measures

Smoothness,
enrichments

GFEM/XFEM versus C^k -GFEM



- Regularity;
- Polynomial reproducibility;
- Efficient enrichment patterns;
- Flat-top property;
- Integration cost;
- Integration of C^0 and C^k -GFEM.



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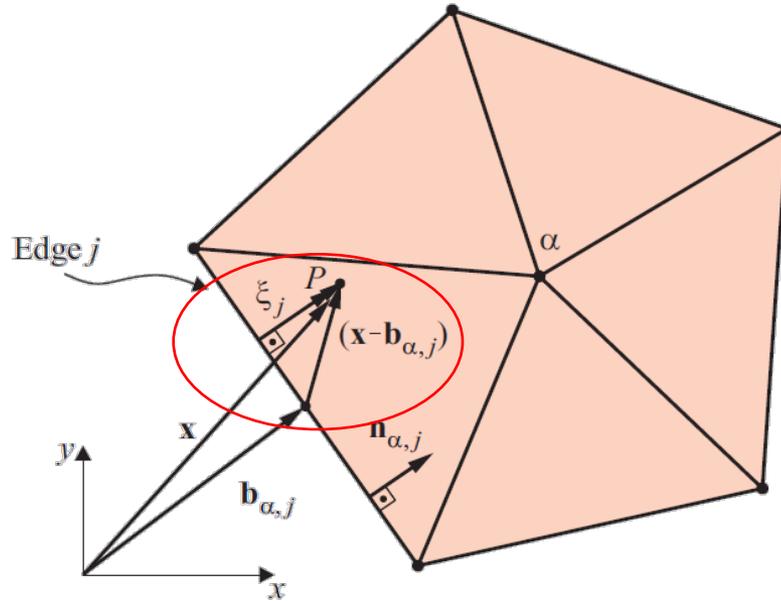
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Concluding remarks

C^k partition of unity – convex clouds



Arbitrary patch shape;
Arbitrary element shape;
free of coordinate mapping;

$$\xi_j(\mathbf{x}) = \mathbf{n}_{\alpha,j} \cdot (\mathbf{x} - \mathbf{b}_{\alpha,j})$$

$$\varepsilon_{\alpha,j}[\xi_j(\mathbf{x})] = \widehat{\varepsilon}_{\alpha,j}(\mathbf{x}) := \begin{cases} e^{-\xi_j^{-\gamma}} & \text{if } 0 < \xi_j \\ 0, & \text{otherwise} \end{cases} \quad k = \infty$$

$$\varepsilon_{\alpha,j}[\xi_j(\mathbf{x})] = \widehat{\varepsilon}_{\alpha,j}(\mathbf{x}) := \begin{cases} (\xi_j/h_j)^p & \text{if } 0 < \xi_j \\ 0, & \text{otherwise} \end{cases} \quad k = p-1$$

Edwards, *C^∞ finite element basis functions*, Report 45,
Institute for Computational Engineering and Sciences – The University of Texas at Austin, 1996

Duarte, Kim and Quaresma, *Arbitrarily smooth generalized finite element approximations*. Computer
Methods in Applied Mechanics and Engineering, 196 (2006)

Barcellos, Mendonça and Duarte, *A C^k continuous generalized finite element formulation applied to
laminated Kirchhoff plate model*. Computational Mechanics, 44 (2009)



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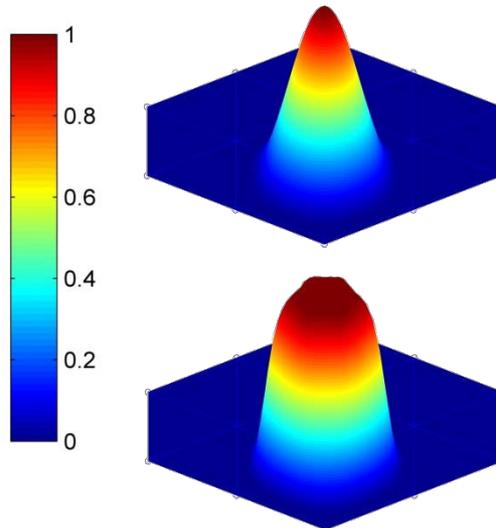
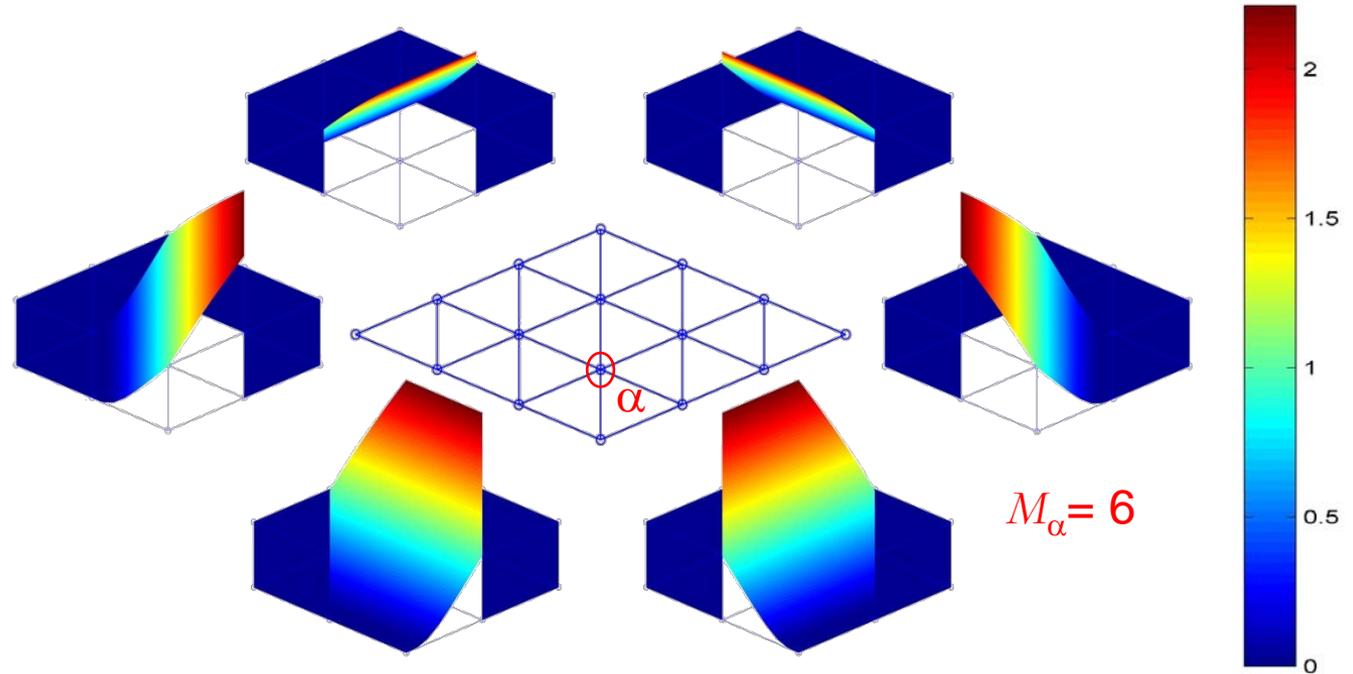
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Some improvements
beyond...

C^∞ partition of unity – convex clouds



$$\mathcal{W}_\alpha(\mathbf{x}) := \prod_{j=1}^{M_\alpha} \varepsilon_{\alpha,j}(\xi_j)$$

$$\varphi_\alpha(\mathbf{x}) = \frac{\mathcal{W}_\alpha(\mathbf{x})}{\sum_{\beta(\mathbf{x})} \mathcal{W}_\beta(\mathbf{x})},$$

$$\beta(\mathbf{x}) \in \{\gamma \mid \mathcal{W}_\gamma(\mathbf{x}) \neq 0\}$$



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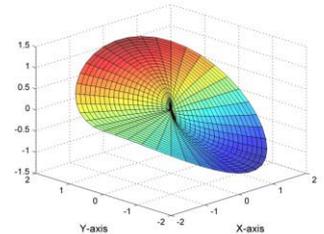
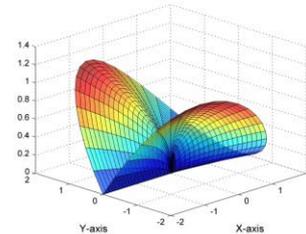
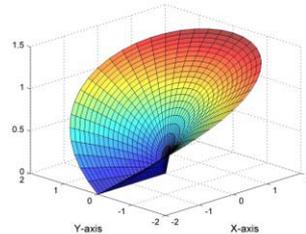
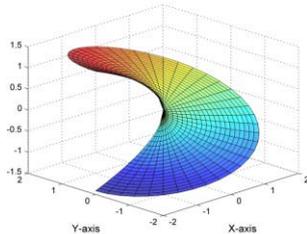
Galerkin approximation

$$\mathbf{u}_p(\mathbf{x}) = \sum_{\alpha=1}^N \varphi_{\alpha}(\mathbf{x}) \left\{ u_{\alpha} + \sum_{i=1}^{q_{\alpha}} \mathcal{L}_{\alpha i}(\mathbf{x}) b_{\alpha i} + \sum_{j=1}^{q_{\alpha}^s} \mathcal{L}_{\alpha j}^s b_{\alpha j}^s \right\}$$

if $p=3$ $\mathcal{L}_{\alpha 9}(x, y) = \left\{ \bar{x}, \bar{y}, \bar{x}^2, \bar{x} \bar{y}, \bar{y}^2, \bar{x}^3, \bar{x}^2 \bar{y}, \bar{x} \bar{y}^2, \bar{y}^3 \right\}$

e.g. $\bar{x} := \frac{(x - x_{\alpha})}{h_{\alpha}}$ for reducing mesh dependences

$$q_{\alpha}^s = 4 \text{ or } 0$$



$$\mathcal{L}_{\alpha 4}^s(r, \theta) = \left\{ \sqrt{r} \sin\left(\frac{\theta}{2}\right), \sqrt{r} \cos\left(\frac{\theta}{2}\right), \sqrt{r} \sin\left(\frac{\theta}{2}\right) \sin(\theta), \sqrt{r} \cos\left(\frac{\theta}{2}\right) \sin(\theta) \right\}$$

Belytschko and Black, *Elastic crack growth in finite elements with minimal remeshing*. International Journal for Numerical Methods in Engineering, 45 (1999)



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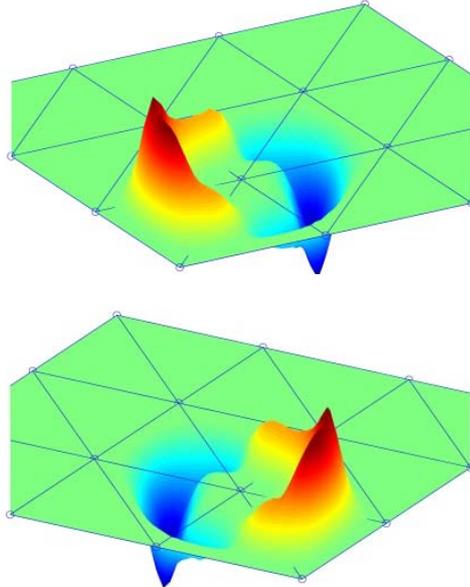
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Polynomial reproducibility of the approximation



$b=p+1$ for C^0 PoU (conventional tent
FEM shape function)

$b=p$ for C^k PoU

p = degree of polynomial enrichment

Mendonça, Barcellos and Torres, **Robust C^k/C^0 generalized FEM approximations for higher-order conformity requirements: application to Reddy's HSDT model for anisotropic laminated plates.** Composite Structures, 96 (2013)

Mendonça, Barcellos and Torres, **Analysis of anisotropic Mindlin plate model by continuous and non-continuous GFEM.** Finite Element in Analysis and Design, 47 (2011)

Barcellos, Mendonça and Duarte, **A C^k continuous generalized finite element formulation applied to laminated Kirchhoff plate model.** Computational Mechanics, 44 (2009)



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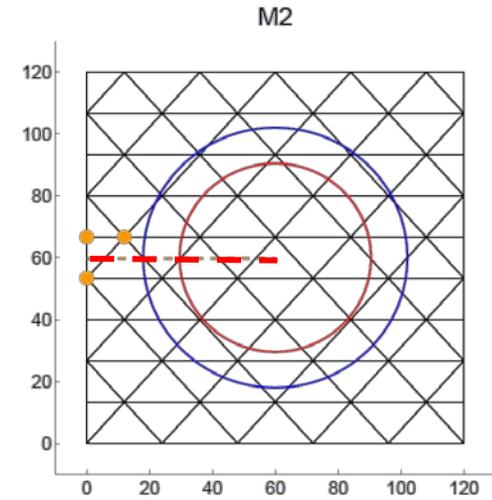
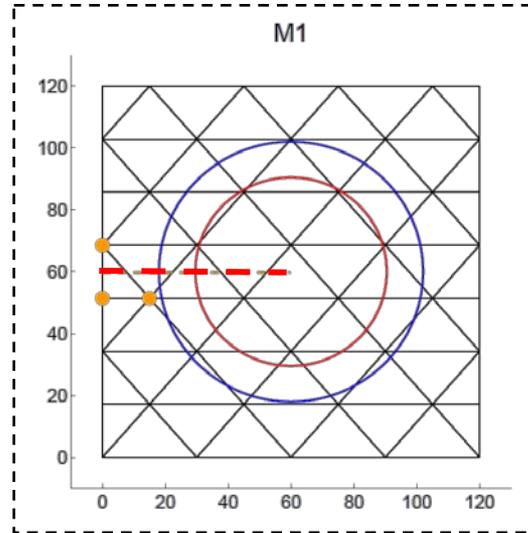
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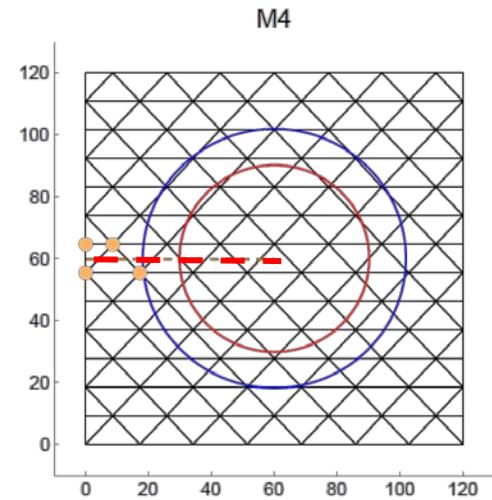
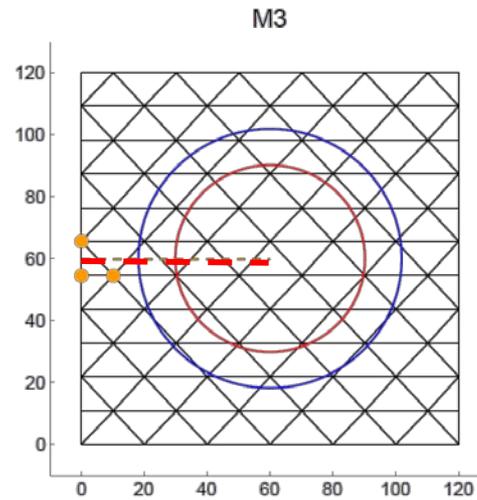
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Geometric enrichment pattern



$R1 > R2$



branch functions on orange nodes and inside the circles,
and uniform polynomial enrichment



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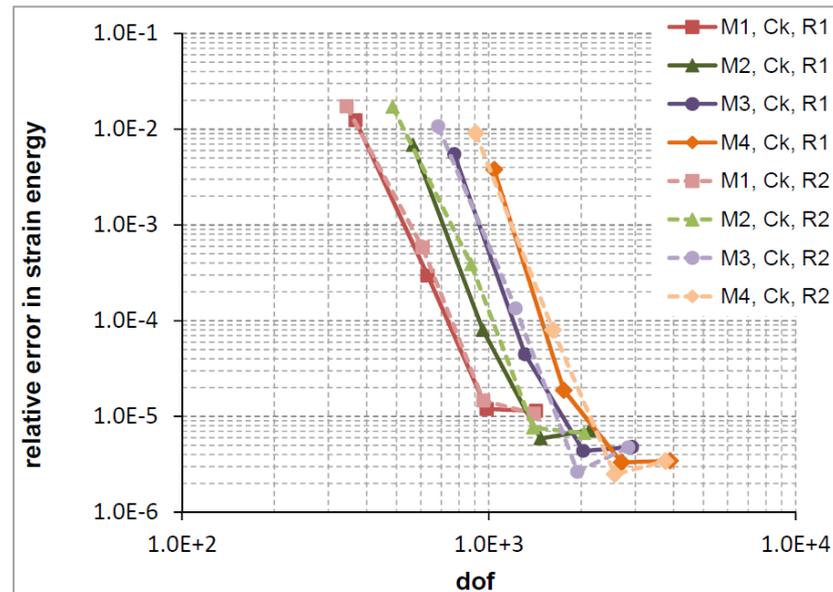
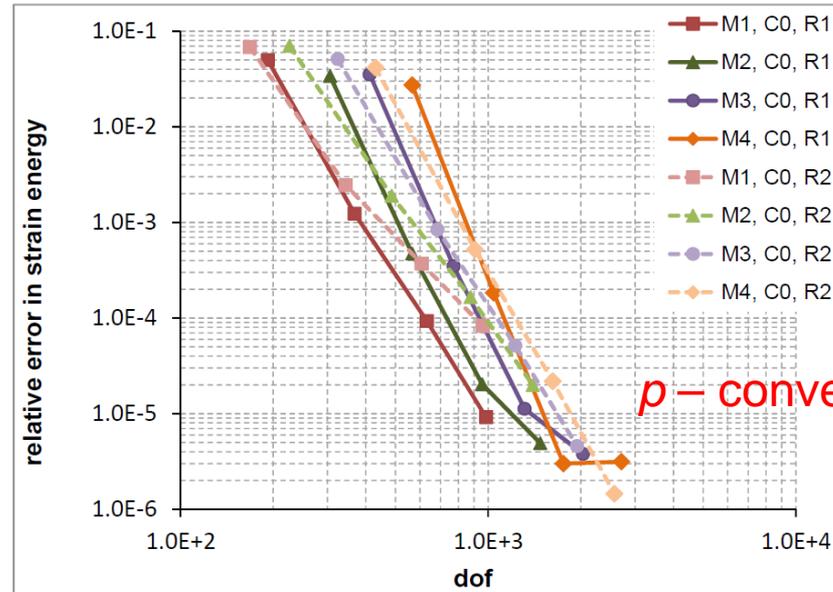
Convergence in terms of global values

mode I opening crack
loading

geometric pattern of
singular enrichment

+

uniform polynomial
enrichment





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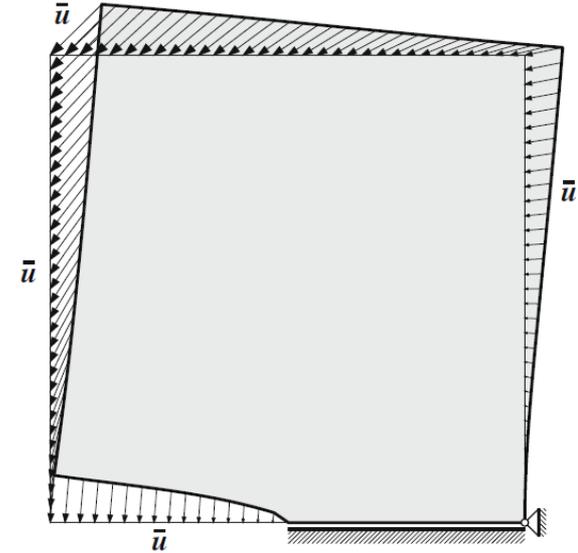
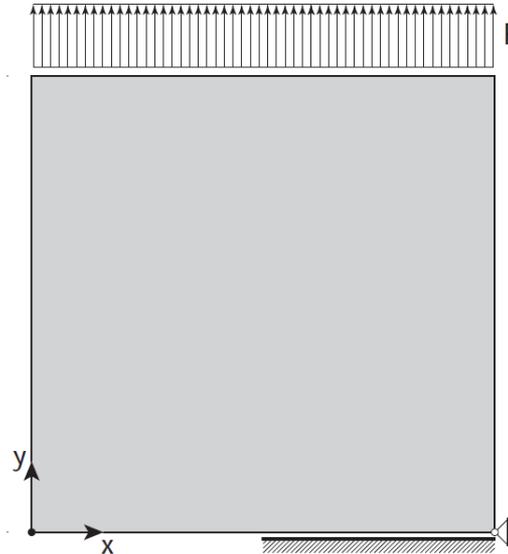
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Configurational mechanics



$$\mathbb{L}^T \boldsymbol{\sigma}(u) + \mathbf{b} = \mathbf{0} \quad , \text{ on } \Omega$$

$$\mathbb{L}^T \boldsymbol{\Sigma}(u) + \boldsymbol{\rho} = \mathbf{0} \quad , \text{ on } \Omega$$

$$\boldsymbol{\sigma} = \{ \sigma_x, \sigma_y, \tau_{xy} \}^T \quad \mathbf{b} = \{ b_x, b_y \}^T \quad \boldsymbol{\Sigma} = \{ \Sigma_x, \Sigma_y, \Sigma_{xy}, \Sigma_{yx} \}^T \quad \boldsymbol{\rho} = \{ \rho_x, \rho_y \}^T$$

Eshelby, *The force on an elastic singularity*. Philosophical Transactions of the Royal Society A: mathematical, physical and engineering sciences, 244 (1951)

Kienzler and Herrmann, *Mechanics in material space with applications to defect and fracture mechanics*. Springer, 2000

Ruter and Stein, *On the duality of global finite element discretization error control in small strain Newtonian and Eshelbian mechanics*. Technische Mechanik, 23 (2003)



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Variational balance of material linear momentum

strong-form

$$\Sigma = \{\Sigma_x, \Sigma_y, \Sigma_{xy}, \Sigma_{yx}\}^T$$

$$\mathbb{L}^T \Sigma(\mathbf{u}) + \boldsymbol{\rho} = \mathbf{0} \quad , \text{ where}$$

Eshelby tensor

$$\boldsymbol{\rho} = \{\rho_x, \rho_y\}^T$$

inhomogeneity force

$$\mathbb{L} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & 0 \\ 0 & \frac{\partial}{\partial x} \end{bmatrix}$$

and defining

$$\Sigma(\mathbf{u}) = \mathfrak{W}(\mathbf{u}) \mathbb{I} - \mathbb{L}^T(\mathbf{u}) \boldsymbol{\sigma}(\mathbf{u})$$

$$\mathfrak{W} = \frac{1}{2} \boldsymbol{\sigma}^T \boldsymbol{\varepsilon} \quad \mathbb{I}^T = \{1, 1, 0, 0\}$$

$$\underline{\mathbb{L}}(\mathbf{u}) = \underline{\mathbb{L}} \underline{\mathbf{I}} \mathbf{u} \quad \underline{\mathbb{L}} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 & \frac{\partial}{\partial y} & 0 \\ 0 & \frac{\partial}{\partial y} & 0 & \frac{\partial}{\partial x} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix} \quad \underline{\mathbf{I}} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \\ 1 & 0 \\ 0 & 1 \end{bmatrix}$$

weak-form

$$\int_{\Omega} (\underline{\mathbb{L}}\mathbf{v})^T \Sigma l_z d\Omega = \int_{\Omega} (\mathbf{v})^T \boldsymbol{\rho} l_z d\Omega \quad \mathbf{v} = \{v_x, v_y\}^T$$



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Post-processing of nodal configurational forces

$$\int_{\Omega} (\mathbb{L}\mathbf{v})^T \boldsymbol{\Sigma} l_z d\Omega = \int_{\Omega} (\mathbf{v})^T \boldsymbol{\rho} l_z d\Omega, \quad \forall \mathbf{v} \mid \mathbf{v} = \mathbf{0} \text{ on } \partial\Omega$$

Eshelby tensor $\boldsymbol{\Sigma}(\mathbf{u}) = \mathfrak{W}(\mathbf{u}) \mathbb{I} - \underline{\mathbb{L}}^T(\mathbf{u}) \boldsymbol{\sigma}(\mathbf{u})$

$$\therefore \sum_{\alpha=1}^N \mathbf{U}_{\alpha}^T \left[- \int \int_{\Omega^e} \mathbb{B}_{\alpha}^T \boldsymbol{\Sigma} l_z dx dy + \int \int_{\Omega^e} \hat{\boldsymbol{\varphi}}_{\alpha}^T \boldsymbol{\rho}_{\alpha} l_z dx dy \right] = 0$$

where $\mathbb{B}_{\alpha} = \mathbb{L} \hat{\boldsymbol{\varphi}}_{\alpha}$

thus, defining

$$\mathbf{G}_{\alpha}^e = \left\{ \begin{array}{c} G_{x\alpha}^e \\ G_{y\alpha}^e \end{array} \right\} = \int \int_{\Omega^e} \hat{\boldsymbol{\varphi}}_{\alpha}^T \boldsymbol{\rho}_{\alpha} l_z dx dy = \int \int_{\Omega^e} \mathbb{B}_{\alpha}^T \boldsymbol{\Sigma} l_z dx dy$$

$$\mathbf{G}_{\alpha} = \bigcup_{e=1}^{N_{ad}} \mathbf{G}_{\alpha}^e$$

nodal
configurational
force



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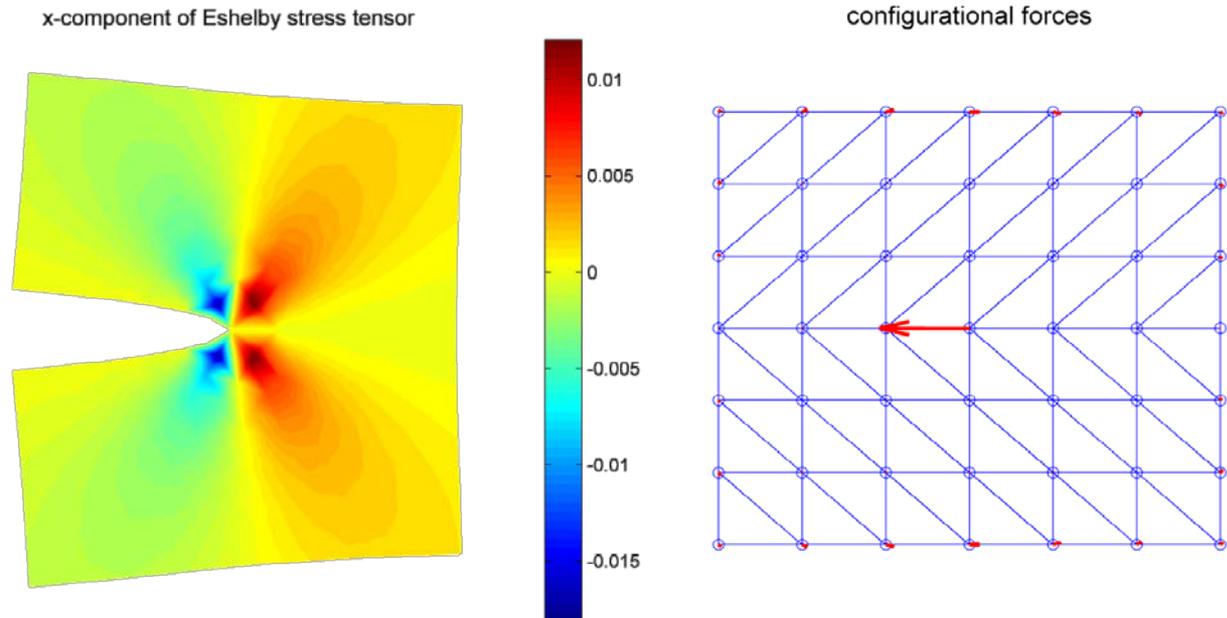
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Local measure using configurational forces



Eshelby, *The force on an elastic singularity*. Philosophical Transactions of the Royal Society A: mathematical, physical and engineering sciences, 244 (1951)

Kienzler and Herrmann, *Mechanics in material space with applications to defect and fracture mechanics*. Springer, 2000

Mueller and Maugin, *On material forces and finite element discretizations*. Computational Mechanics, 29 (2002)

Glaser and Steinmann, *On material forces within the extended finite element method*. Proceedings of the sixth European Solid Mechanics Conference (2006)

Häusler, Lindhorst and Horst, *Combination of the material force concept and the extended finite element method for mixed mode crack growth simulations*. International Journal for Numerical Methods in Engineering, 85 (2011)



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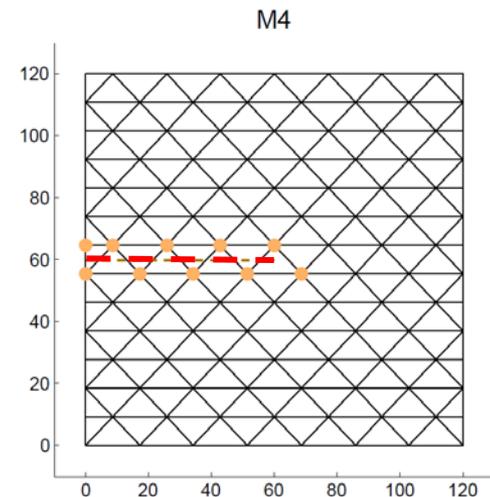
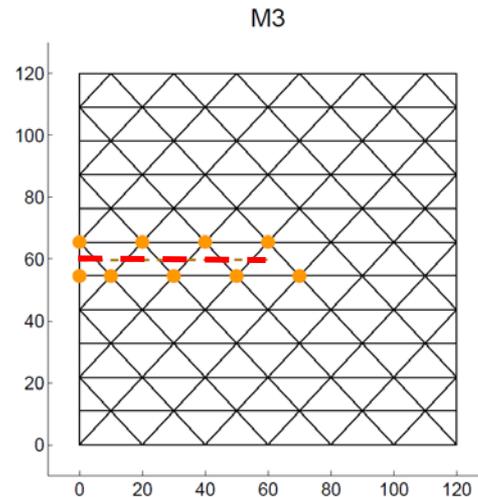
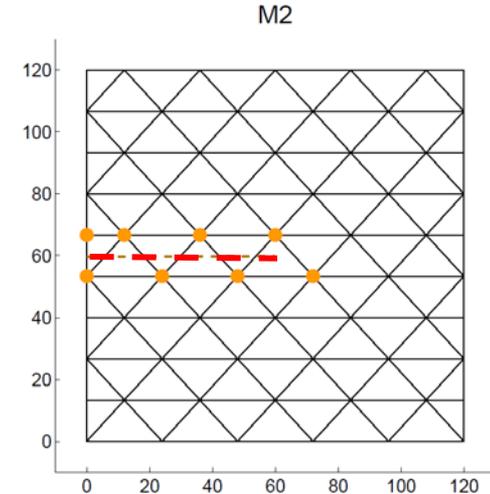
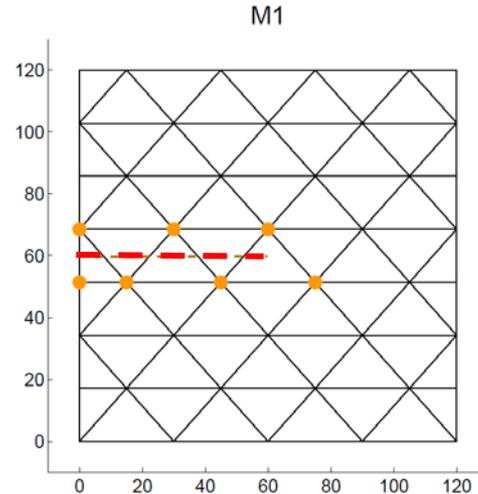
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Topologic enrichment pattern



- Branch functions on orange nodes;
- Uniform p -enrichment
- Local p -enrichment around the crack tip



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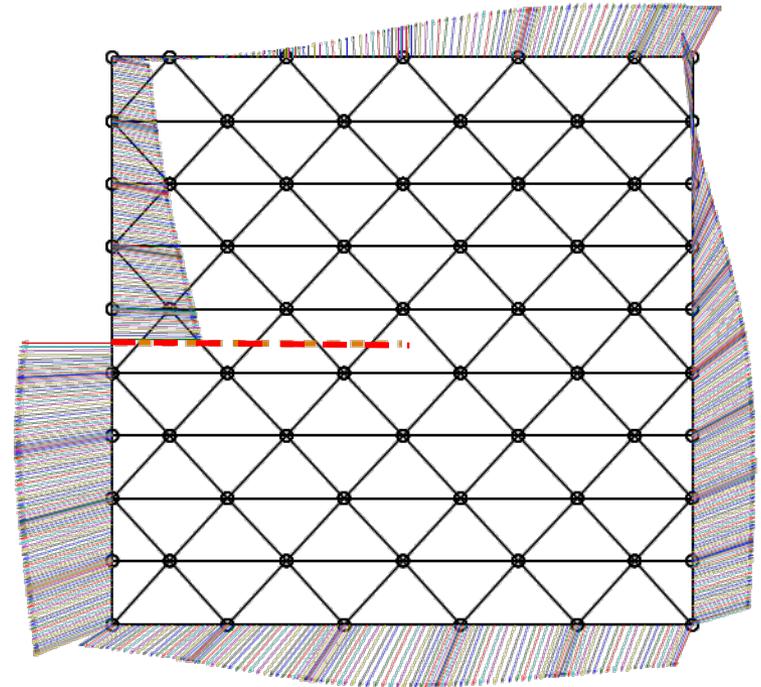
Mixed mode loading

$$K_I = 1.0, K_{II} = 1.0$$

$$\sigma_x = \frac{K_I}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \left[1 - \text{sen}\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{3\theta}{2}\right) \right] - \frac{K_{II}}{\sqrt{2\pi r}} \text{sen}\left(\frac{\theta}{2}\right) \left[2 + \cos\left(\frac{\theta}{2}\right) \cos\left(\frac{3\theta}{2}\right) \right]$$

$$\sigma_y = \frac{K_I}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \left[1 + \text{sen}\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{3\theta}{2}\right) \right] + \frac{K_{II}}{\sqrt{2\pi r}} \text{sen}\left(\frac{\theta}{2}\right) \cos\left(\frac{\theta}{2}\right) \cos\left(\frac{3\theta}{2}\right)$$

$$\tau_{xy} = \frac{K_I}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{\theta}{2}\right) \cos\left(\frac{3\theta}{2}\right) + \frac{K_{II}}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \left[1 - \text{sen}\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{3\theta}{2}\right) \right]$$



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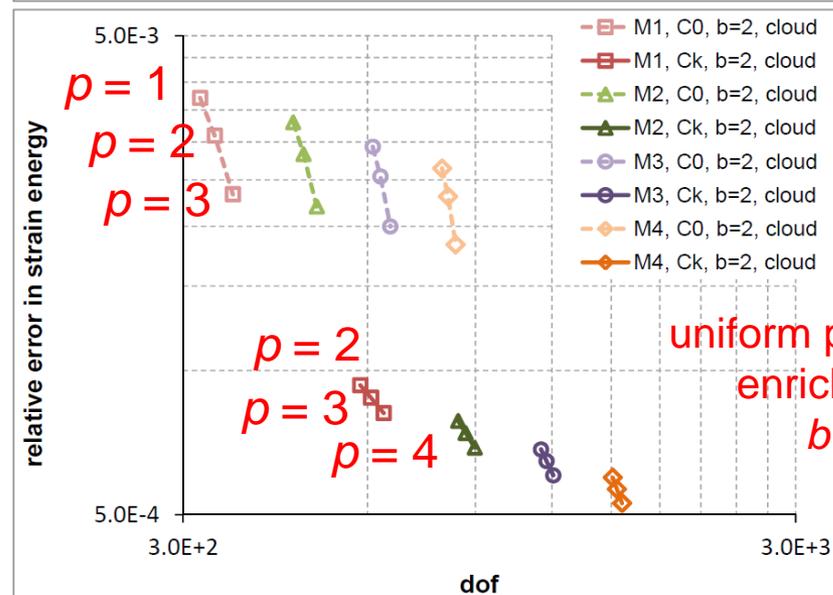
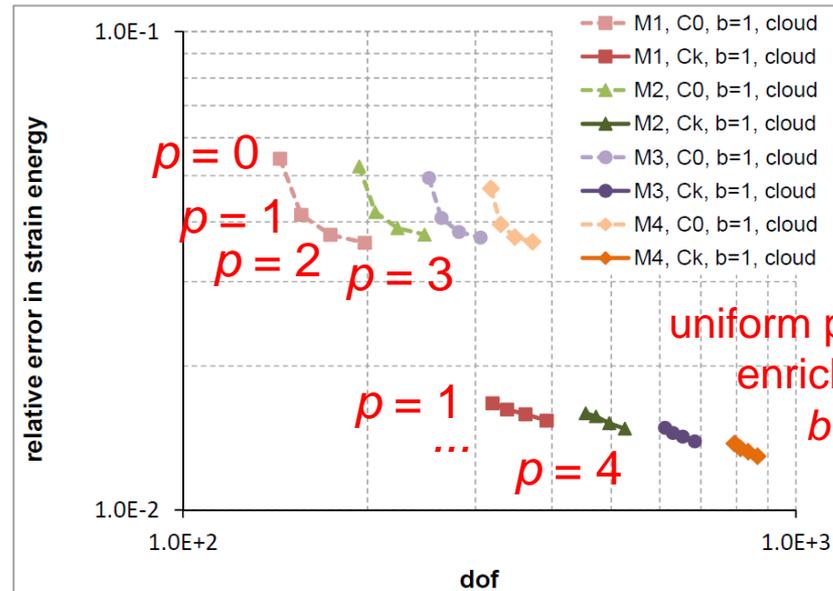
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Convergence in global values

Topologic pattern of
singular enrichment

+

Localized polynomial
enrichment, p





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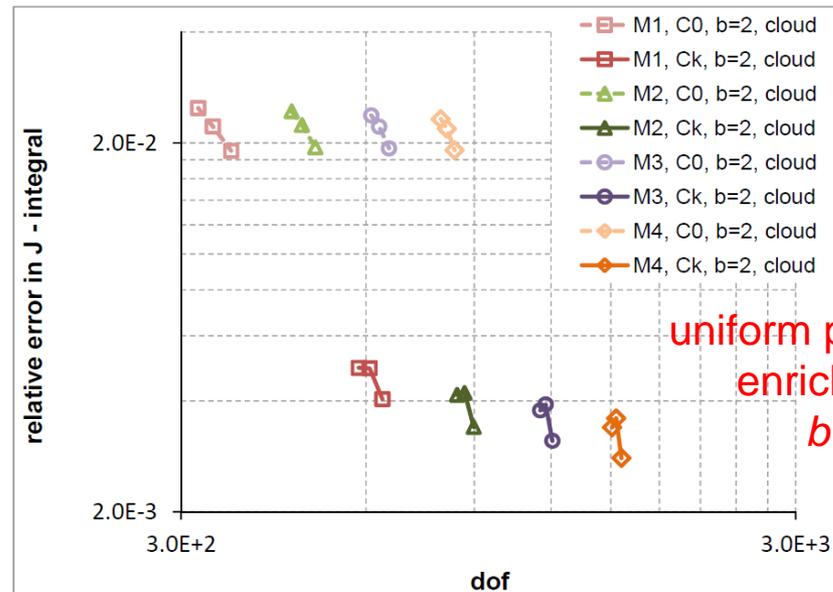
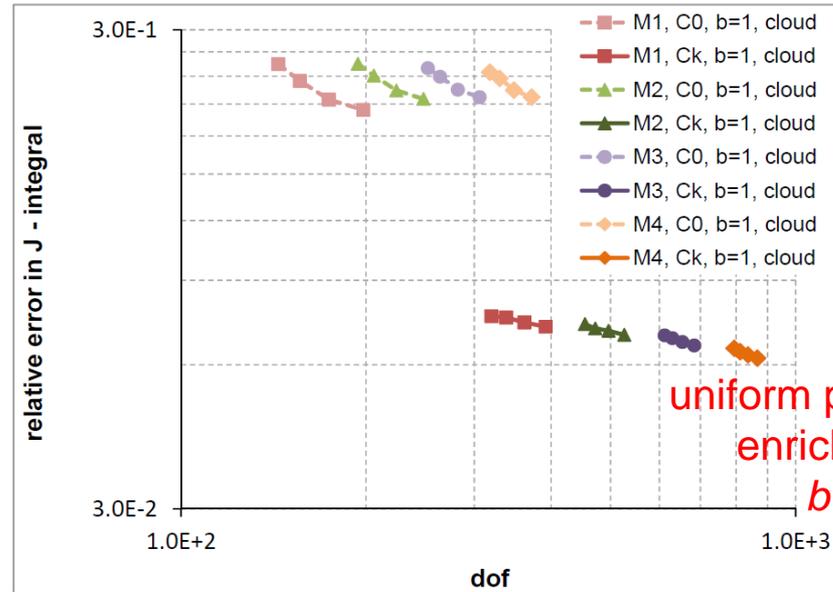
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Localized polynomial
enrichment

$$\mathcal{R} = (K_I^2 + K_{II}^2) / E^*$$

$$E^* = E(1 - \nu^2)$$

Convergence of J-integral





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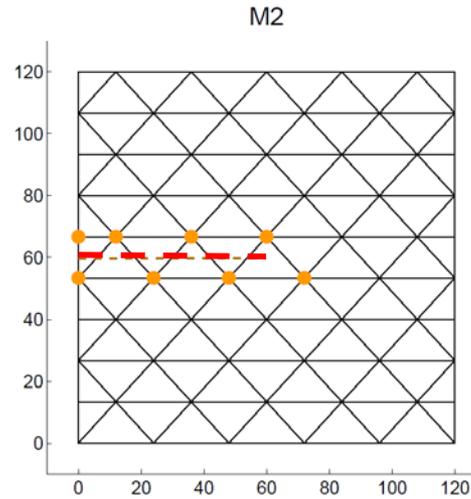
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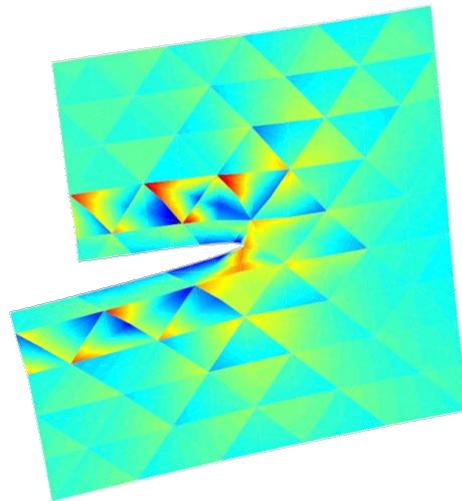
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Exact error dispersion: γ -stress

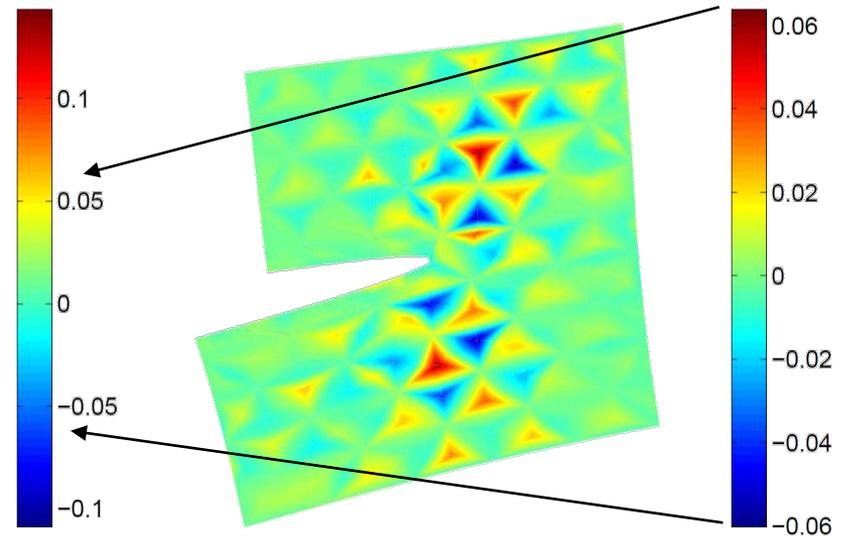
Topologic pattern of
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Uniform polynomial
enrichment,
 $b = 1$



C_0 , 191 DOFs



C_∞ , 451 DOFs



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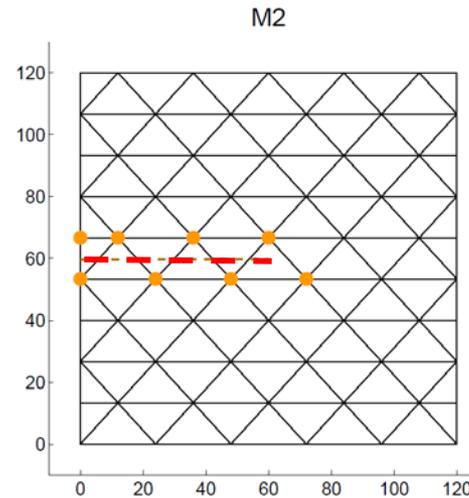
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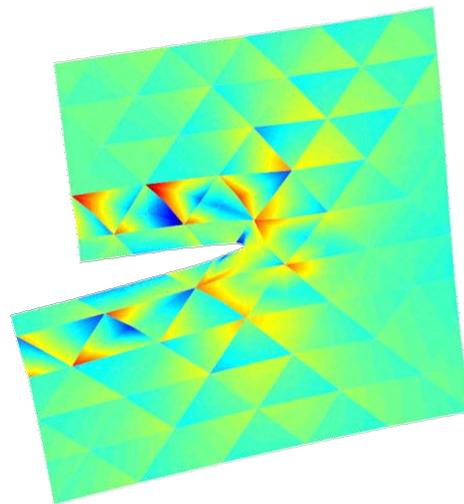
topologic pattern of
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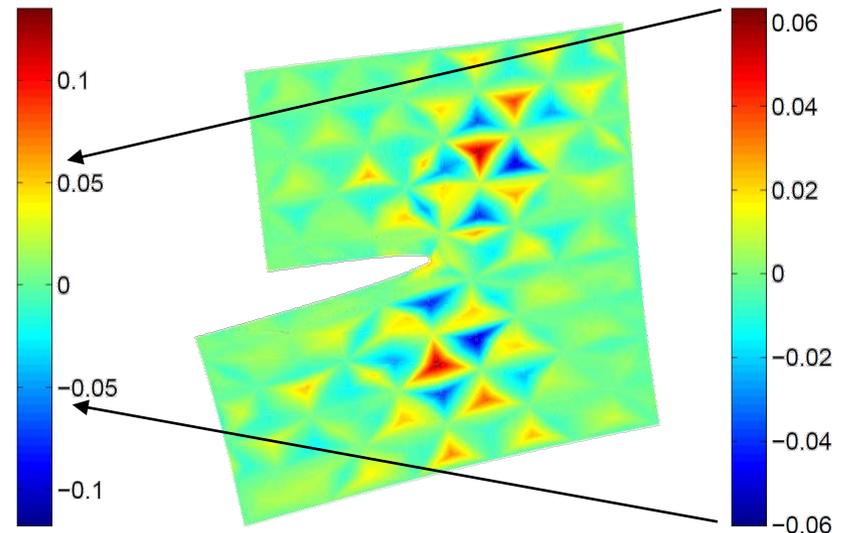
uniform polynomial
enrichment,
 $b = 1$

+

localized polynomial
enrichment,
 $p = 2$



C_0 , 203 DOFs



C_∞ , 469 DOFs



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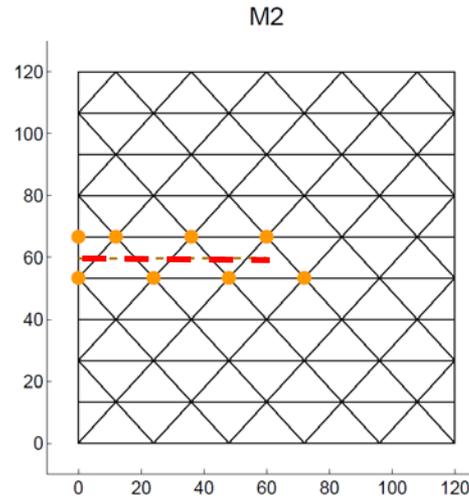
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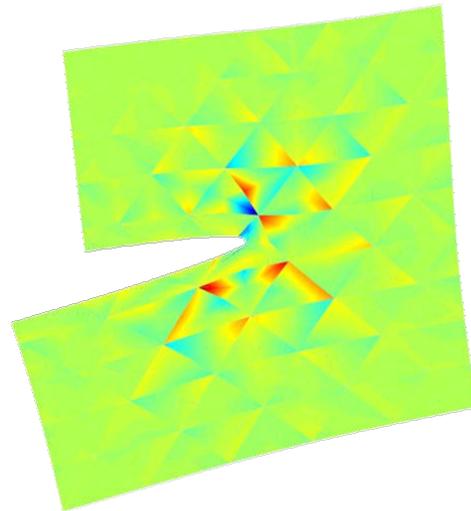
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Exact error dispersion: γ -stress

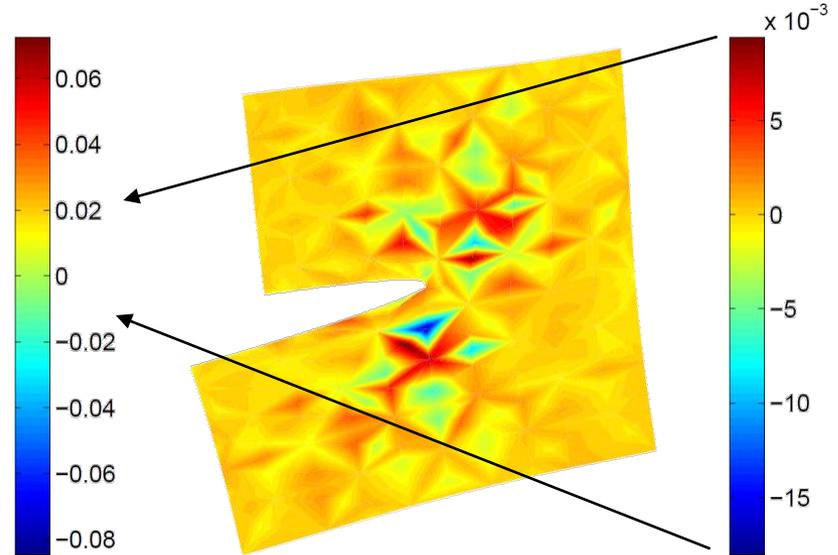
topologic pattern of
singular enrichment



uniform polynomial
enrichment,
 $b = 2$



C_0 , 451 DOFs



C_∞ , 841 DOFs



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Continuous partition of
unity with C^k -GFEM

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subspace

Quality assessment
through global
measures

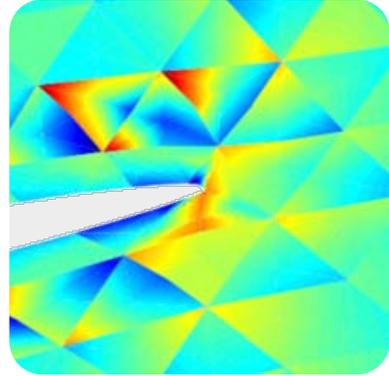
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Smoothness,
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Concluding remarks

Collecting data of the three cases...

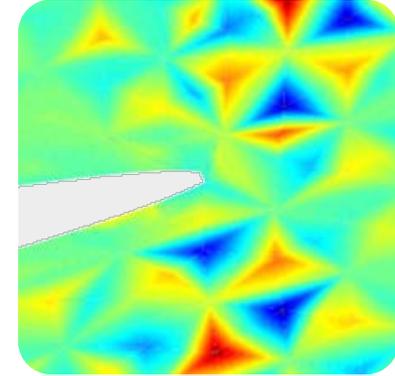


C^0

$b=1$, 191 DOFs, 145×10^{-3}

$b=1$, $p=2$, 203 DOFs, 130×10^{-3}

$b=2$, 451 DOFs, -85×10^{-3}



C^∞

$b=1$, 451 DOFs, 62×10^{-3}

$b=1$, $p=2$, 469 DOFs, 62×10^{-3}

$b=2$, 841 DOFs, -19×10^{-3}



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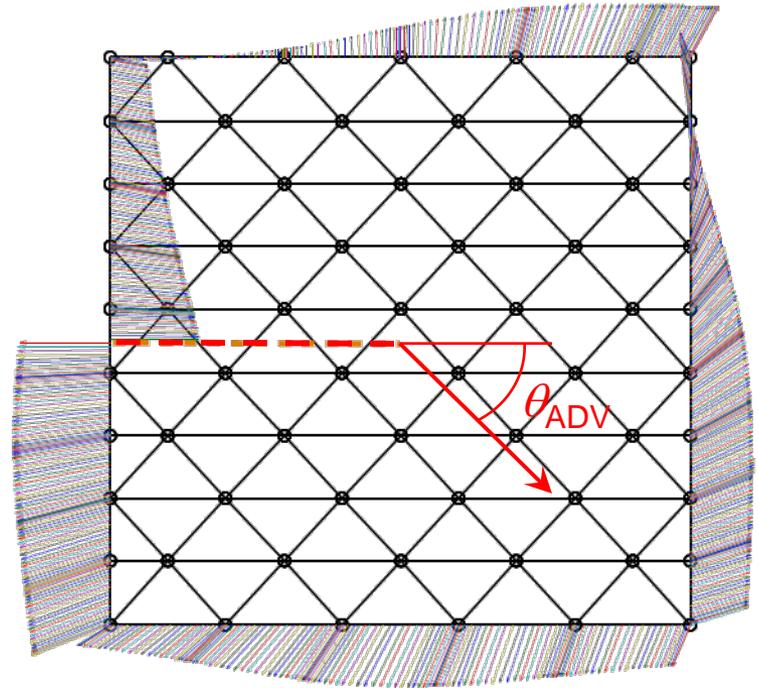
Angle of crack growth: θ_{ADV}

$$K_I = 1.0, K_{II} = 1.0$$

$$\sigma_x = \frac{K_I}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \left[1 - \text{sen}\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{3\theta}{2}\right) \right] - \frac{K_{II}}{\sqrt{2\pi r}} \text{sen}\left(\frac{\theta}{2}\right) \left[2 + \cos\left(\frac{\theta}{2}\right) \cos\left(\frac{3\theta}{2}\right) \right]$$

$$\sigma_y = \frac{K_I}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \left[1 + \text{sen}\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{3\theta}{2}\right) \right] + \frac{K_{II}}{\sqrt{2\pi r}} \text{sen}\left(\frac{\theta}{2}\right) \cos\left(\frac{\theta}{2}\right) \cos\left(\frac{3\theta}{2}\right)$$

$$\tau_{xy} = \frac{K_I}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{\theta}{2}\right) \cos\left(\frac{3\theta}{2}\right) + \frac{K_{II}}{\sqrt{2\pi r}} \cos\left(\frac{\theta}{2}\right) \left[1 - \text{sen}\left(\frac{\theta}{2}\right) \text{sen}\left(\frac{3\theta}{2}\right) \right]$$



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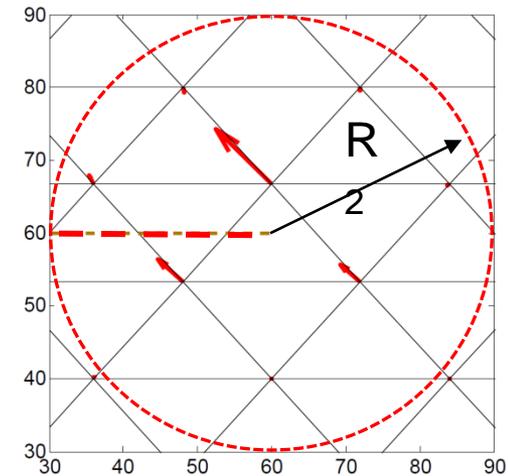
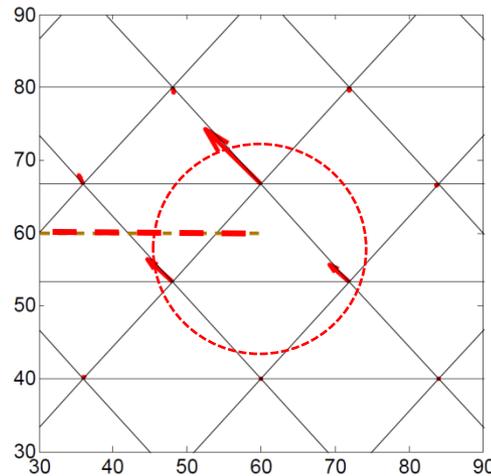
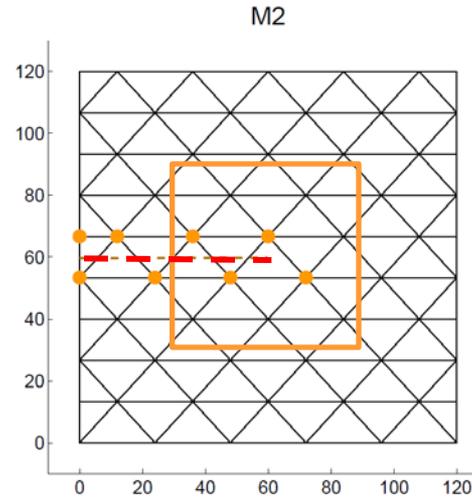
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Angle of crack advance: θ_{ADV}



Configurational forces are restorative forces !



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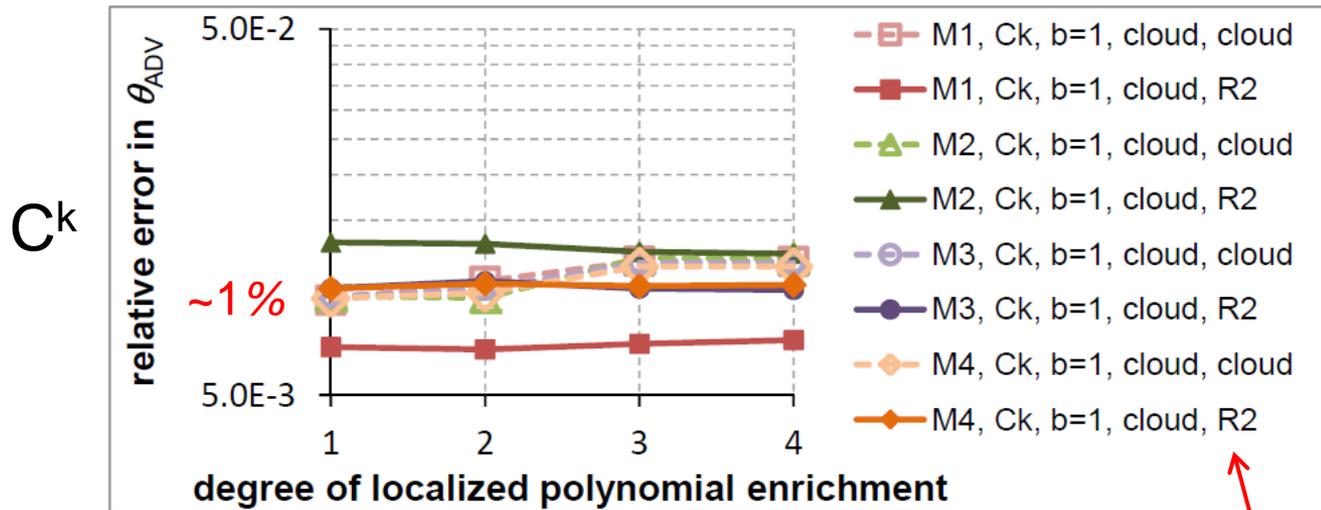
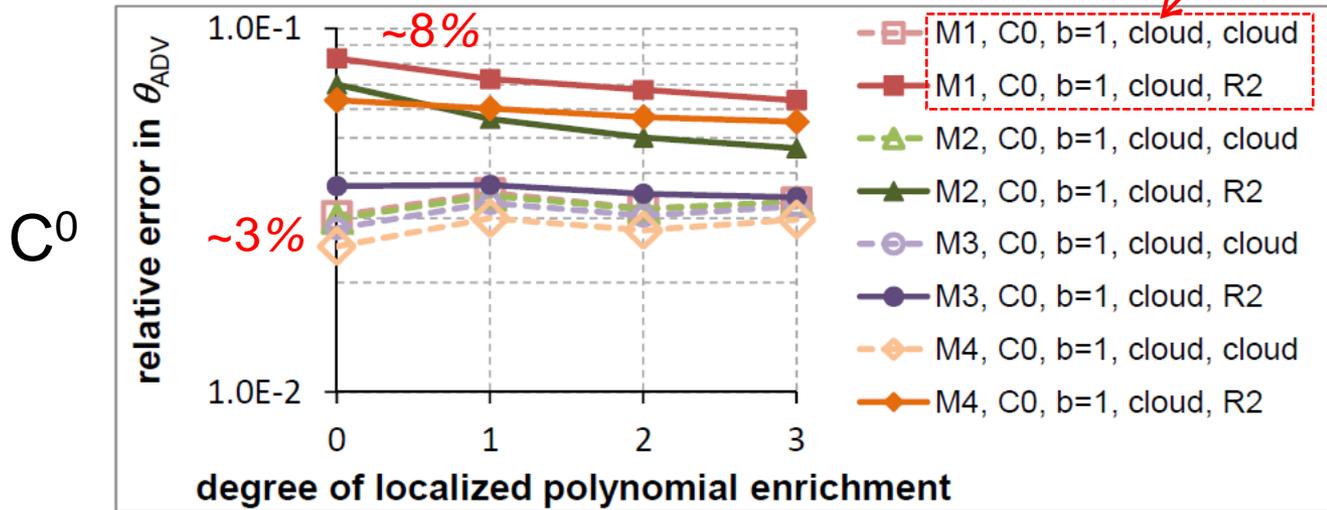
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Angle of crack advance: θ_{ADV} , $b = 1$

Crack tip
enrichment



Topologic pattern of
singular enrichment

Uniform polynomial
enrichment,
 $b = 1$

Nodes for the
configurational force



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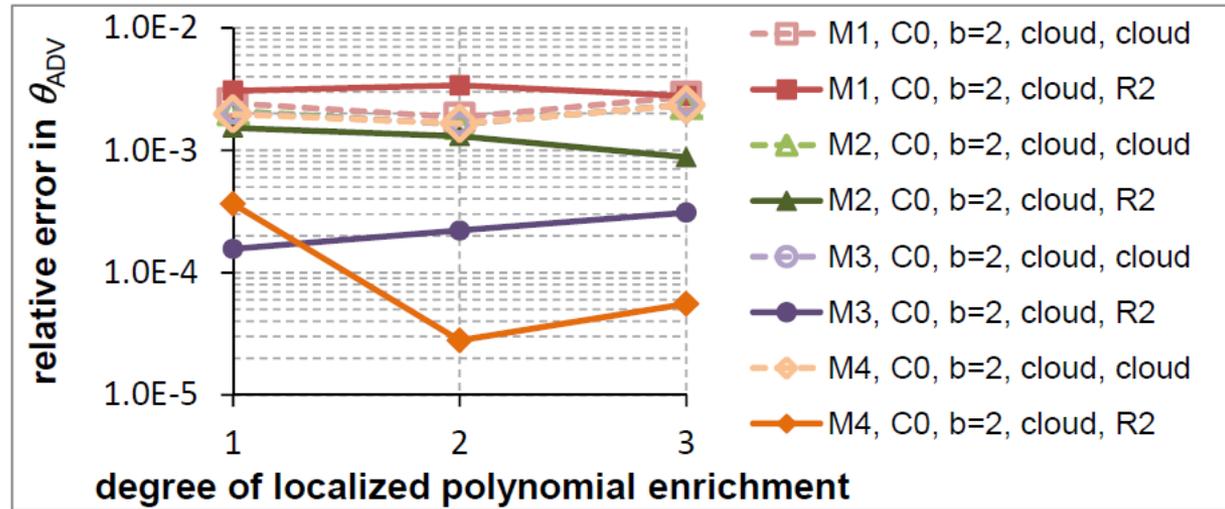
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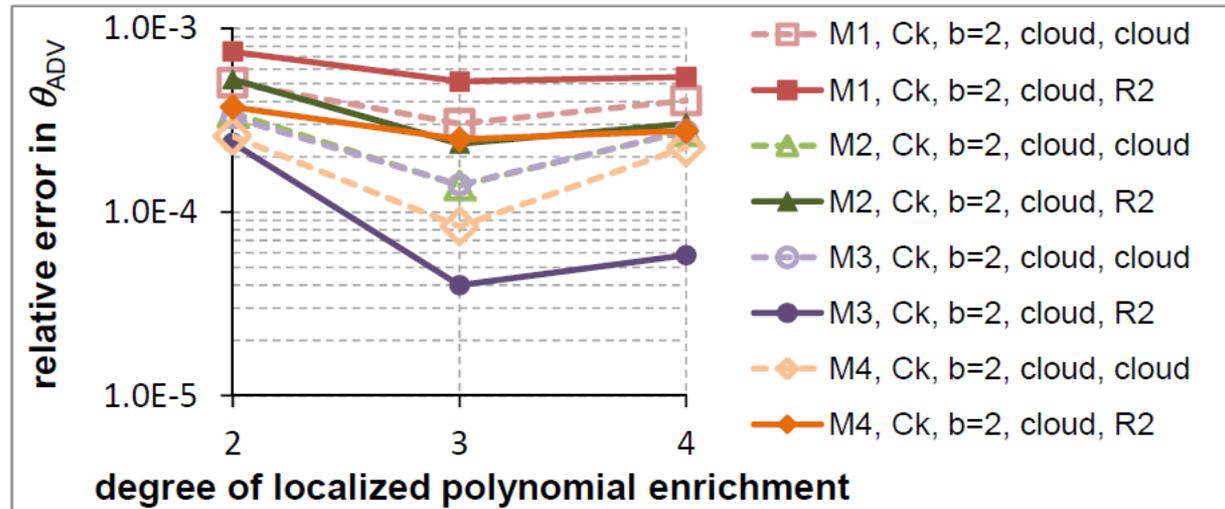
Concluding remarks

Angle of crack advance: θ_{ADV} , $b = 2$

C^0



C^k



Topologic pattern of
singular enrichment

Uniform polynomial
enrichment,
 $b = 2$



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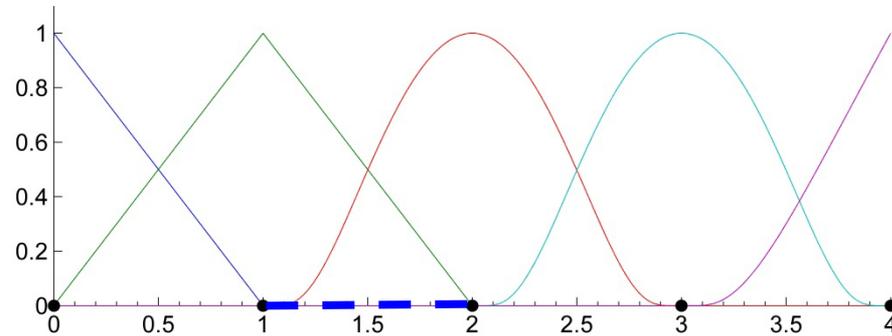
Smoothness,
enrichments and
conditioning

Some improvements

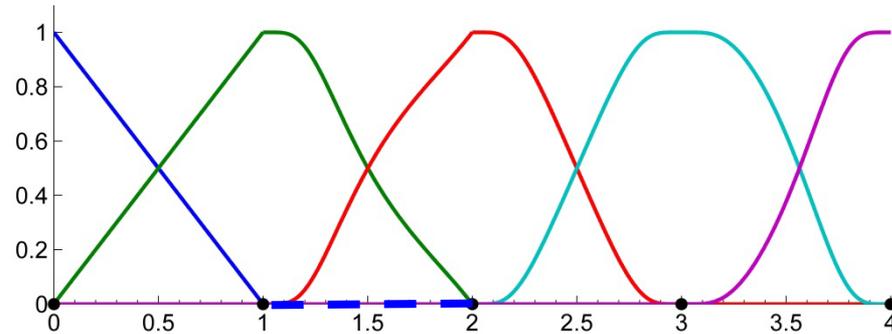
Concluding remarks

Joining C^k -GFEM and C^0 -GFEM

C^0 and C^k weighting functions



Partition of unity



Shepard equation ->

$$\varphi_\alpha(\mathbf{x}) = \frac{W_\alpha(\mathbf{x})}{\sum_{\beta(\mathbf{x})} W_\beta(\mathbf{x})}$$

→ Straightforward!!!

→ Application of smooth PoUs only where they are convenient !

Concluding remarks

C^k -GFEM shows better estimates of J and θ_{ADV} than C^0 -GFEM:

- Probably due to the well determined flat-top;
- Smoothness seems to reduce transition effects around singular enrichment;
- Lower dependence with enrichment pattern;
- Lower dependence on the size of the region used to compute configurational forces;
- In general, C^k -GFEM combines:
 - Higher regularity;
 - Flat-top property;
 - Definition on global coordinates -> admits extremely distorted meshes;
 - Compact support.

Disadvantages: - Lower polynomial reproducibility than C^0 -GFEM /XFEM;
- Higher integration cost;

Solution: - Apply smooth PoUs only at convenient zones of the model !

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Thank you!

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